

Clebsch-Gordan Coefficient Example Phys 402

The eigenfunctions of J^2 can be expressed as linear combinations of states with different values of m_ℓ and m_s using the world-famous Clebsch-Gordan coefficients

$(C_{m_\ell m_s m_j}^{\ell s j})$ as:

$$|j m_j\rangle = \sum_{m_\ell+m_s=m_j} C_{m_\ell m_s m_j}^{\ell s j} |\ell m_\ell\rangle |s m_s\rangle \quad [4.185]$$

Where the ket $|\ell m_\ell\rangle$ represents the spherical harmonics $Y_\ell^{m_\ell}$. The C-G coefficient values are given in Table 4.8 on page 188 of Griffiths. Remember that all of the coefficients should appear under a square root, with the minus sign (if any) out front.

Now for an example of how to construct states that are simultaneous eigenfunctions of L^2 , S^2 , J^2 and J_z . Take the case again of hydrogen with $\ell = 1$ and spin $s = 1/2$. How do we find the state with $j = 3/2$ and $m_j = -1/2$ in terms of the $Y_\ell^{m_\ell}$ and spinors? Look at the $1 \times 1/2$ CG Table on page 188. We are led to this table because we are combining an angular momentum vector with $\ell = 1$ and spin vector with $s = 1/2$.

| | | | | | |
|----------------|--------|--------|--------|--------|--------|
| $1 \times 1/2$ | $3/2$ | | | | |
| | $+3/2$ | $3/2$ | $1/2$ | | |
| | $+1$ | $+1/2$ | 1 | $+1/2$ | $+1/2$ |
| | | $+1$ | $-1/2$ | $1/3$ | $2/3$ |
| | | 0 | $+1/2$ | $2/3$ | $-1/3$ |
| | | | 0 | $-1/2$ | $2/3$ |
| | | | -1 | $+1/2$ | $1/3$ |
| | | | -1 | $-1/2$ | $-3/2$ |
| | | | | | $3/2$ |
| 2×1 | 3 | 3 | 2 | | |
| | $+3$ | 3 | 2 | -1 | $-1/2$ |
| | | | | | 1 |

Now look under the column labeled “ $\begin{matrix} 3/2 \\ -1/2 \end{matrix}$ “. It says:

$$\left| \begin{matrix} 3 \\ 2 \end{matrix} - \begin{matrix} 1 \\ 2 \end{matrix} \right\rangle = \sum_{m_\ell+m_s=-1/2} C_{m_\ell m_s -1/2}^{1 \ 1/2 \ 3/2} \left| \begin{matrix} 1 \\ m_\ell \end{matrix} \right\rangle \left| \begin{matrix} 1/2 \\ m_s \end{matrix} \right\rangle$$

$$\left| \begin{matrix} 3 \\ 2 \end{matrix} - \begin{matrix} 1 \\ 2 \end{matrix} \right\rangle = \sqrt{\frac{2}{3}} \left| \begin{matrix} 1 \\ 0 \end{matrix} \right\rangle \left| \begin{matrix} 1/2 \\ -1/2 \end{matrix} \right\rangle + \sqrt{\frac{1}{3}} \left| \begin{matrix} 1 \\ -1 \end{matrix} \right\rangle \left| \begin{matrix} 1/2 \\ 1/2 \end{matrix} \right\rangle$$

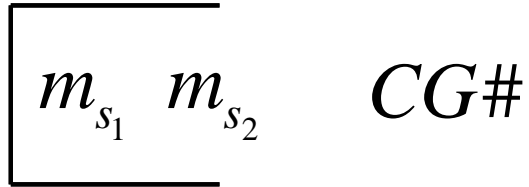
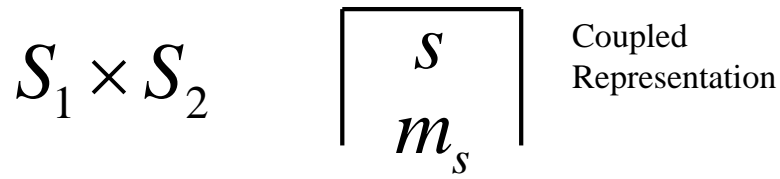
This can be written in a more familiar way in terms of spherical harmonics and spinors as:

$$\left| \begin{matrix} 3 \\ 2 \end{matrix} - \begin{matrix} 1 \\ 2 \end{matrix} \right\rangle = \sqrt{\frac{2}{3}} Y_1^0 \chi_- + \sqrt{\frac{1}{3}} Y_1^{-1} \chi_+$$

One can move back and forth between the coupled and un-coupled representations using the Clebsch-Gordan table on page 188. Here is the schematic layout for the CG table for combining two spins (called \vec{S}_1, \vec{S}_2) to form a total spin $\vec{S} = \vec{S}_1 + \vec{S}_2$ (S^2 has eigenvalue $s(s+1)\hbar^2$):

General Schematic of the C-G Table

$$\vec{S} = \vec{S}_1 + \vec{S}_2$$



Un-Coupled Representation